

# **The Integral Effect & the intrinsic uncertainty (randomness) in dynamical systems**

**Jin-Song von Storch**

**Max-Planck Institute for Meteorology**



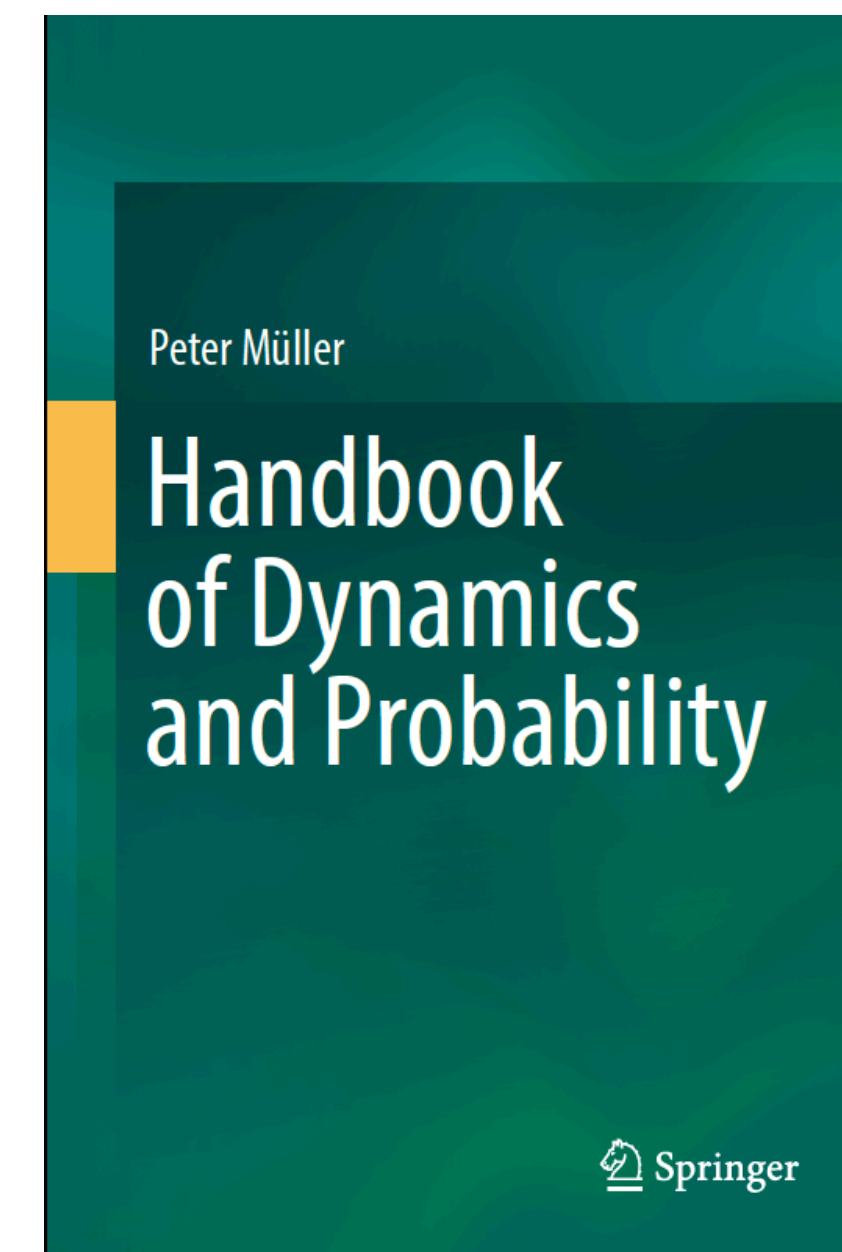
# The Integral Effect & the intrinsic uncertainty (randomness) in dynamical systems

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- Uncertainties arise from
  - lack of knowledge
  - inability of controlling certain things
- Intrinsic uncertainty is a phenomenon which occurs even in case of complete understanding and full control
- Physical origin of intrinsic uncertainty is unknown (apart from that at quantum level)

Peter Müller, 2022



For a dynamical system described by  $\frac{dx}{dt} = f(\mathbf{x})$ , randomness is

- observed in ultra low-frequency variations (ULFV) of equilibrium solution of  $x$

**ULFV not explainable by deterministic working of differential forcing  $f$**

*(On Equilibrium Fluctuations, von Storch, 2022, Tellus)*

- $(2\pi\omega)^2 \Gamma^x(\omega) = \Gamma^f(\omega)$

- At frequency  $\omega = 0$ ,

$$0 \Gamma^x(0) = \Gamma^f(0)$$

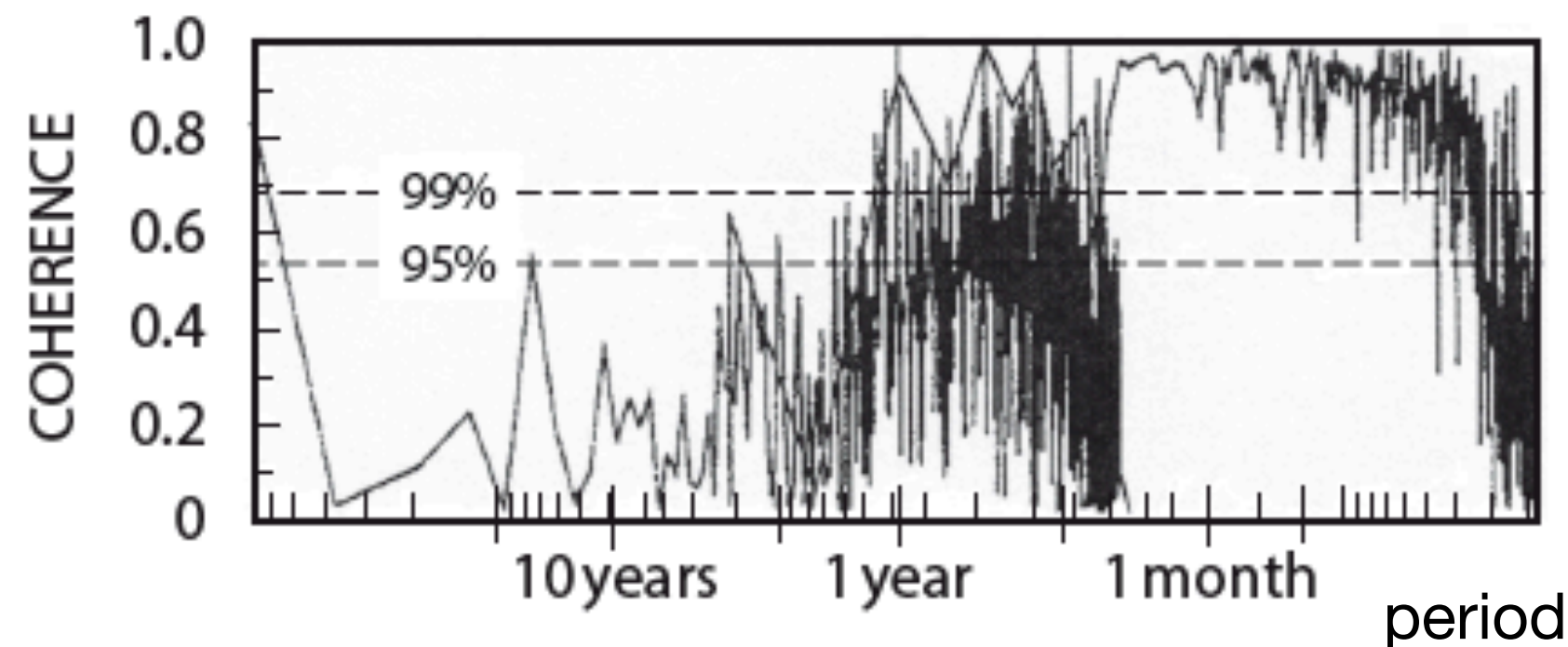
- $\Gamma^f(0)$  must vanish to ensure equilibrium solution

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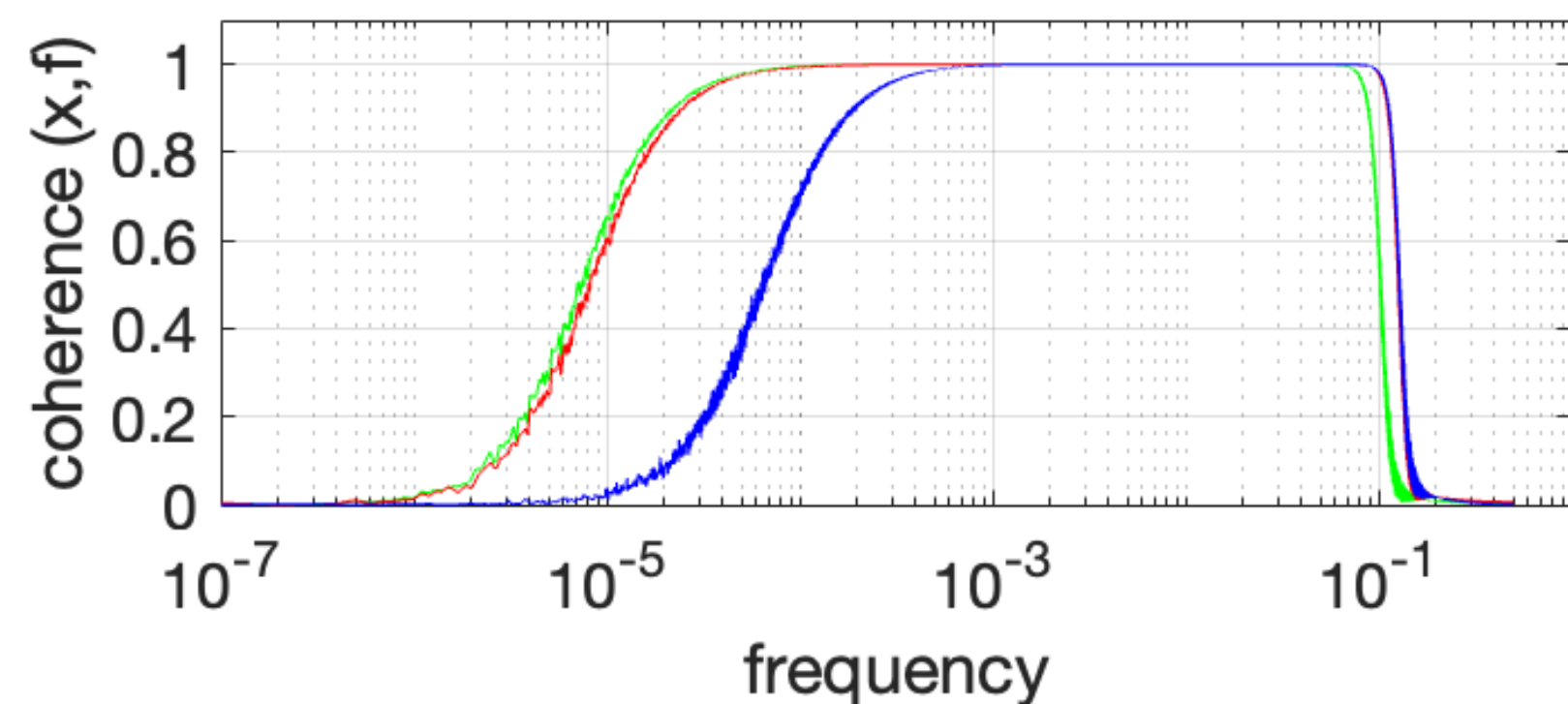
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**ULFV not explainable by deterministic working of differential forcing  $f$**

(On Equilibrium Fluctuations, von Storch, 2022, Tellus)



$x$  = axial angular momentum of the atmosphere simulated by a coupled climate model  
 $f$  = the respective torque  
 (von Storch, 1999)



$x$  = one of the three variables of the Lorenz' 63 model  
 $f$  = the dynamics of the respective variable

**ULFV results from the working of internal forcing  $g_\tau$**

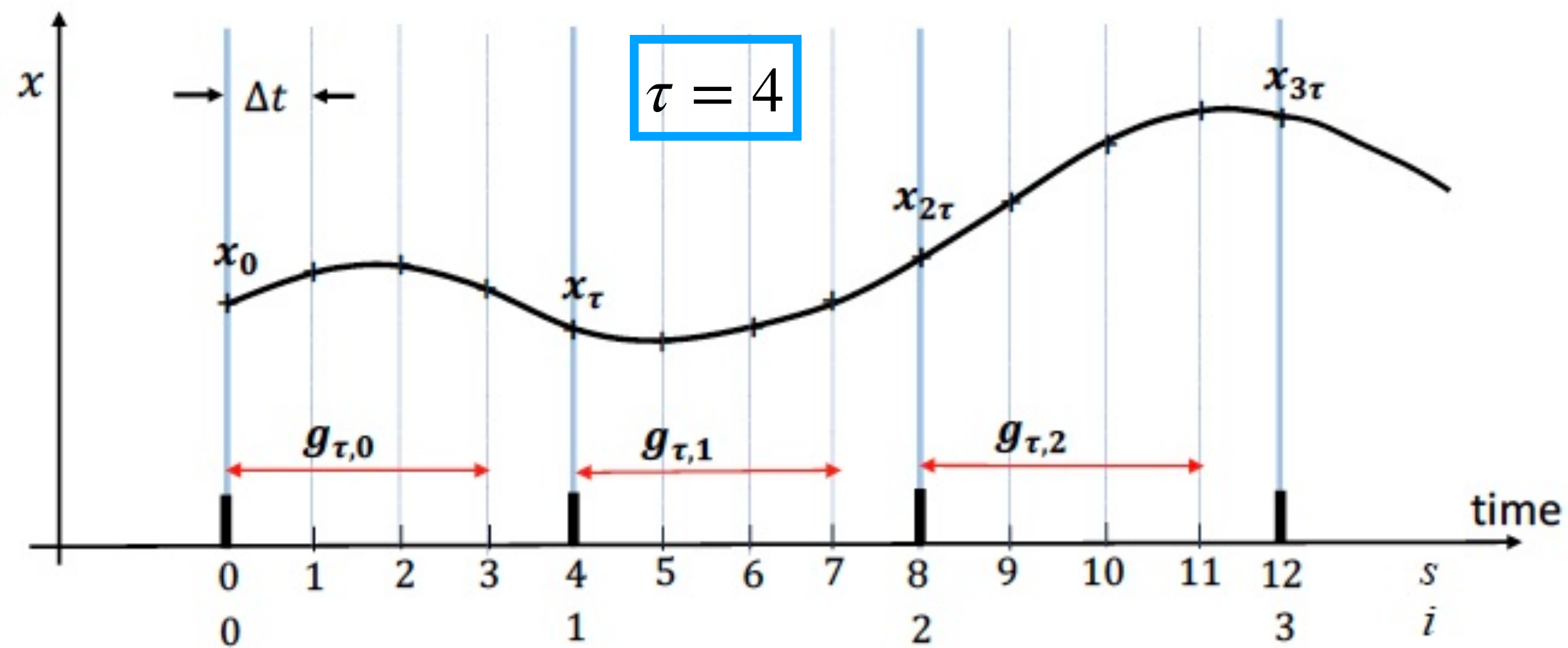
→ The physical origin of randomness

## Preliminaries

- We consider only dynamical systems described by  $dx/dt = f(\mathbf{x})$  that have an equilibrium solution  $x(t)$ , which
  - ◆ varies stationarily for ever when left alone
  - ◆ has a time independent variance - equilibrium variance
- Almost all systems of our interests do not have analytical solutions and have to be solved numerically.
- All numerical evidences are derived from the Lorenz model (1963)

## The integral forcing $g_\tau$

Discretize time axis using increment  $\Delta t$ . Set  $\Delta t = 1$



- $\{x_s\} = \{x_s | s \in \mathbb{Z}_*\}$ : solution at every time steps
- $\{x_{i\tau}\} = \{x_{i\tau} | i \in \mathbb{Z}_*\}$ : solution at every  $\tau$  time steps

Differential forcing  $f_s$ :

- $f_s = f(\mathbf{x}_s)$
- $x_{s+1} = x_s + f_s$

$\tau$ -stepping integral forcing  $g_{\tau,i}$ :

- $g_{\tau,i} = \sum_{s=i\tau}^{(i+1)\tau} f_s$ , for  $\tau \in \mathbb{Z}_+$
- $g_{1,i} = f_i$ , for  $\tau = 1$
- $x_{(i+1)\tau} = x_{i\tau} + g_{\tau,i}$



# Properties of the integral forcing $g_\tau$

An integral forcing can be written as:

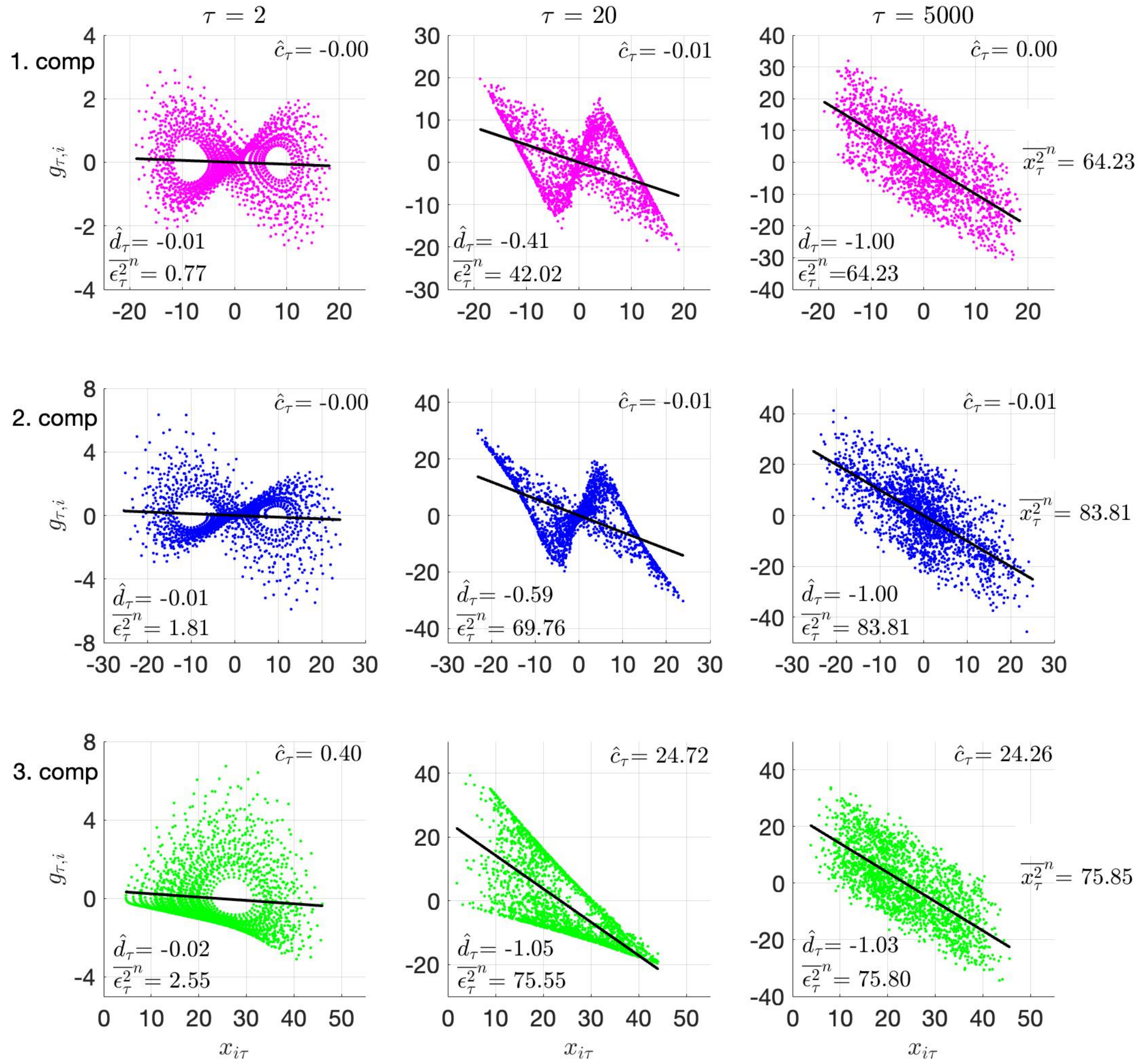
$$g_{\tau,i} = \sum_{s=i\tau}^{(i+1)\tau-1} f_s = \hat{c}_\tau + \hat{d}_\tau x_{i\tau} + \hat{\epsilon}_{\tau,i}$$

$\hat{c}_\tau$ : intercept     $\hat{d}_\tau$ : regression slope

$\hat{\epsilon}_{\tau,i}$ : residual  $g_{\tau,i} - (\hat{c}_\tau + \hat{d}_\tau x_{i\tau})$

$\hat{\cdot}$ : derived from  $n$  data points along an equilibrium solution, here  $n=10^6$

- ◆  $g_{\tau,i}$  becomes increasingly linear in  $x_{i\tau}$  with increasing  $\tau$
- ◆ Once  $g_{\tau,i}$  is linear in  $x_{i\tau}$ ,  $\hat{\epsilon}_{\tau,i}$  behaves like a white noise





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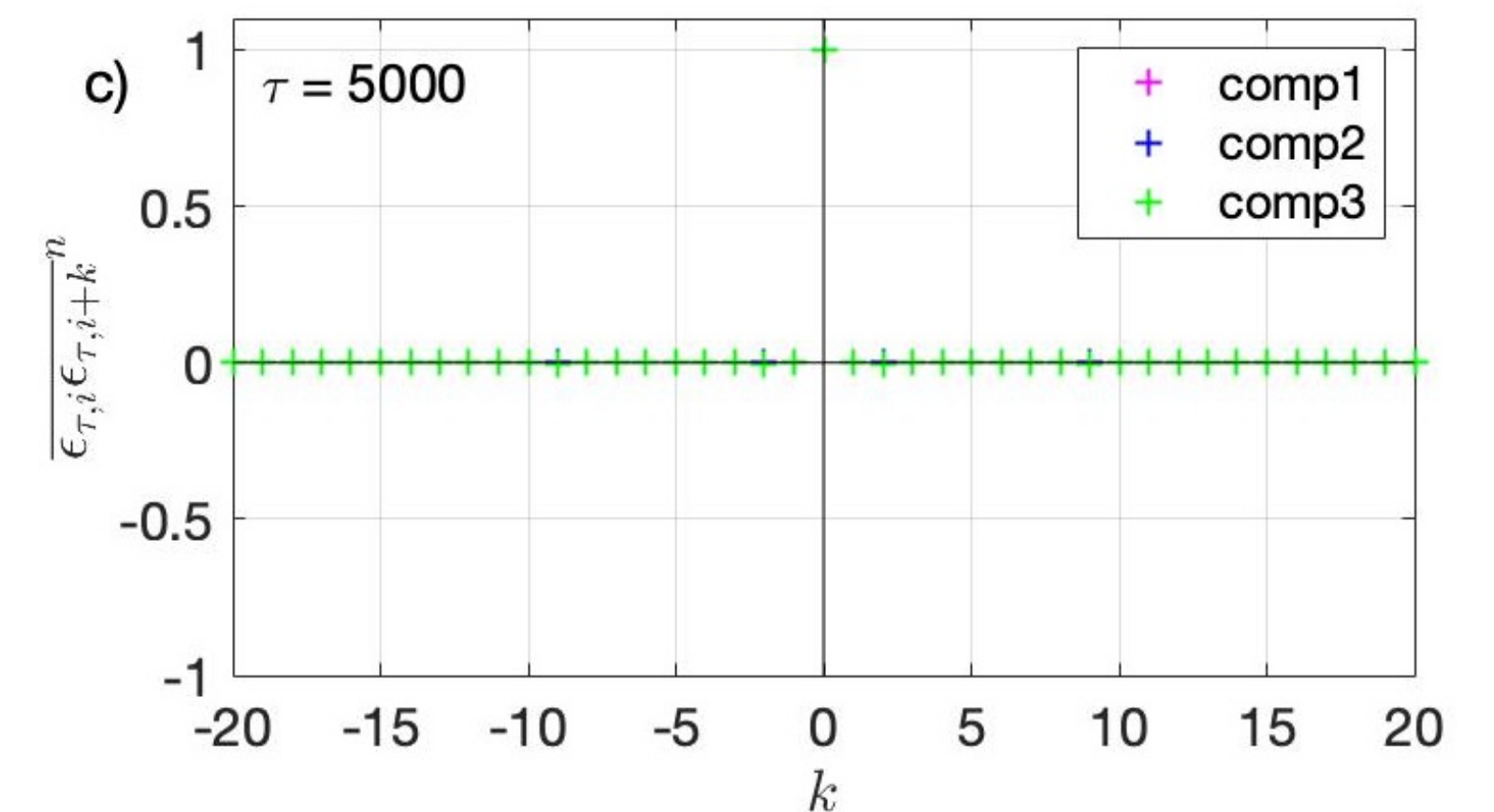
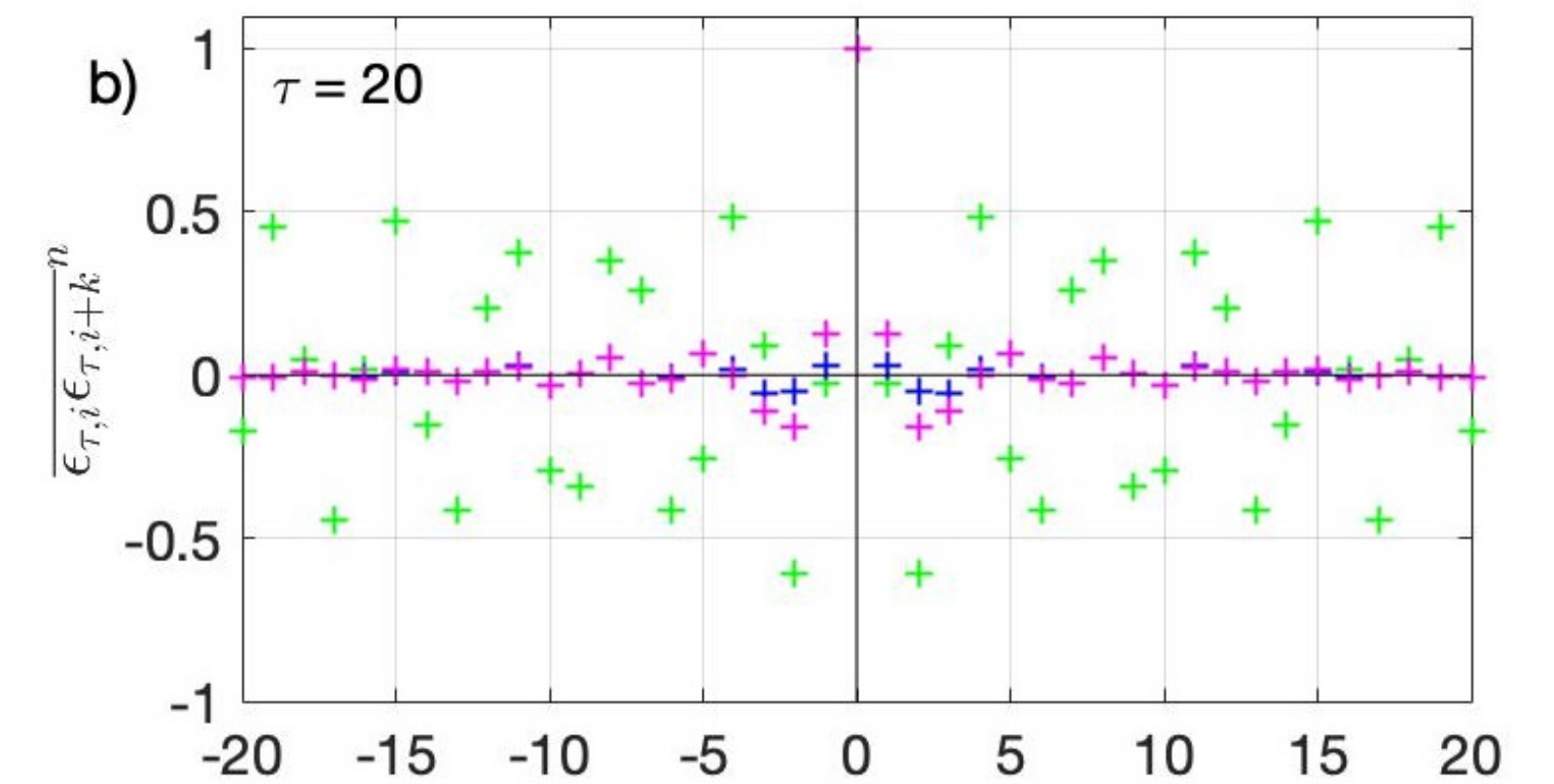
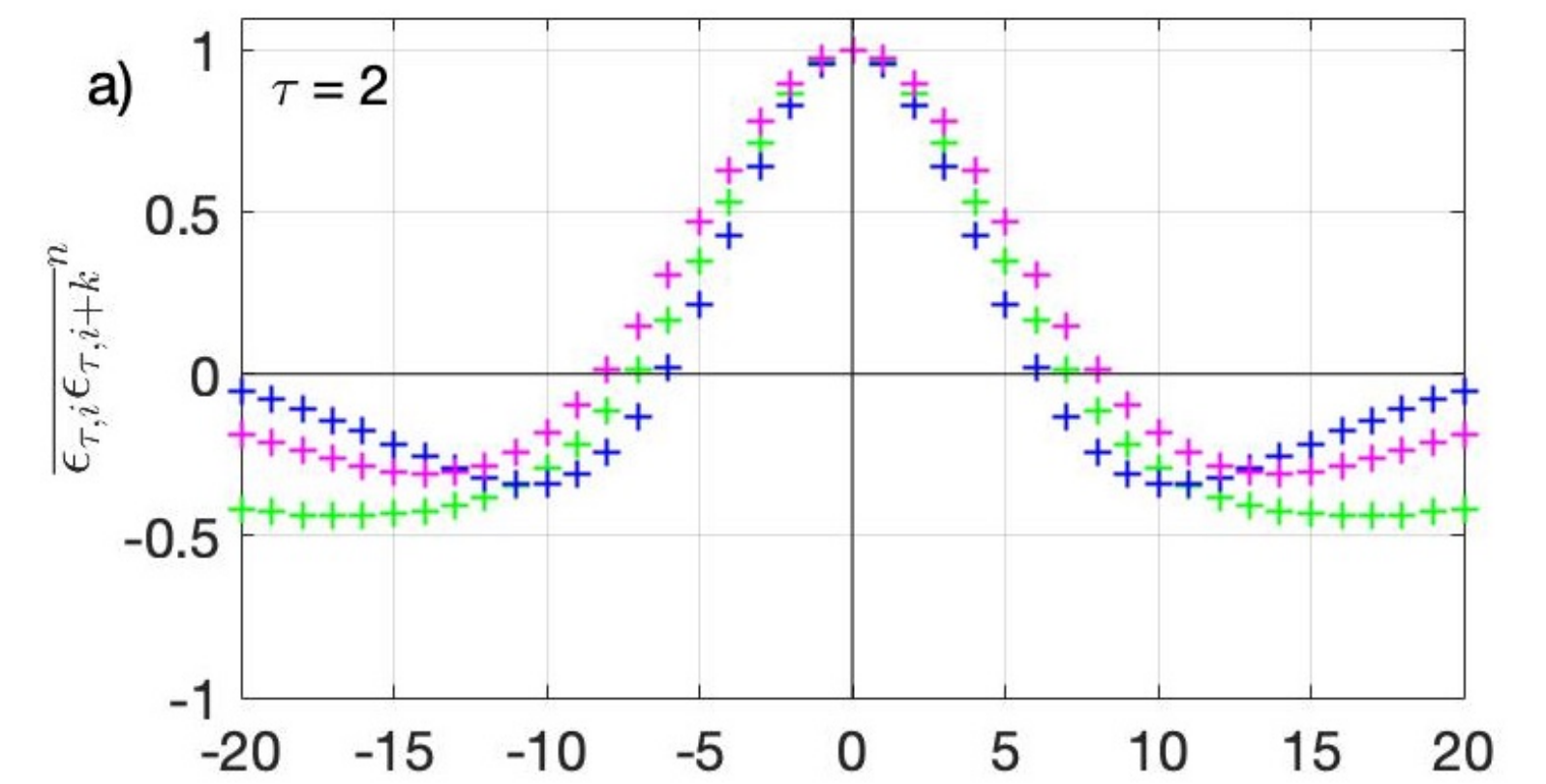
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Auto-correlation function  $\overline{\epsilon_{\tau,i}\epsilon_{\tau,i+k}}^n$  as a function of lag  $k$





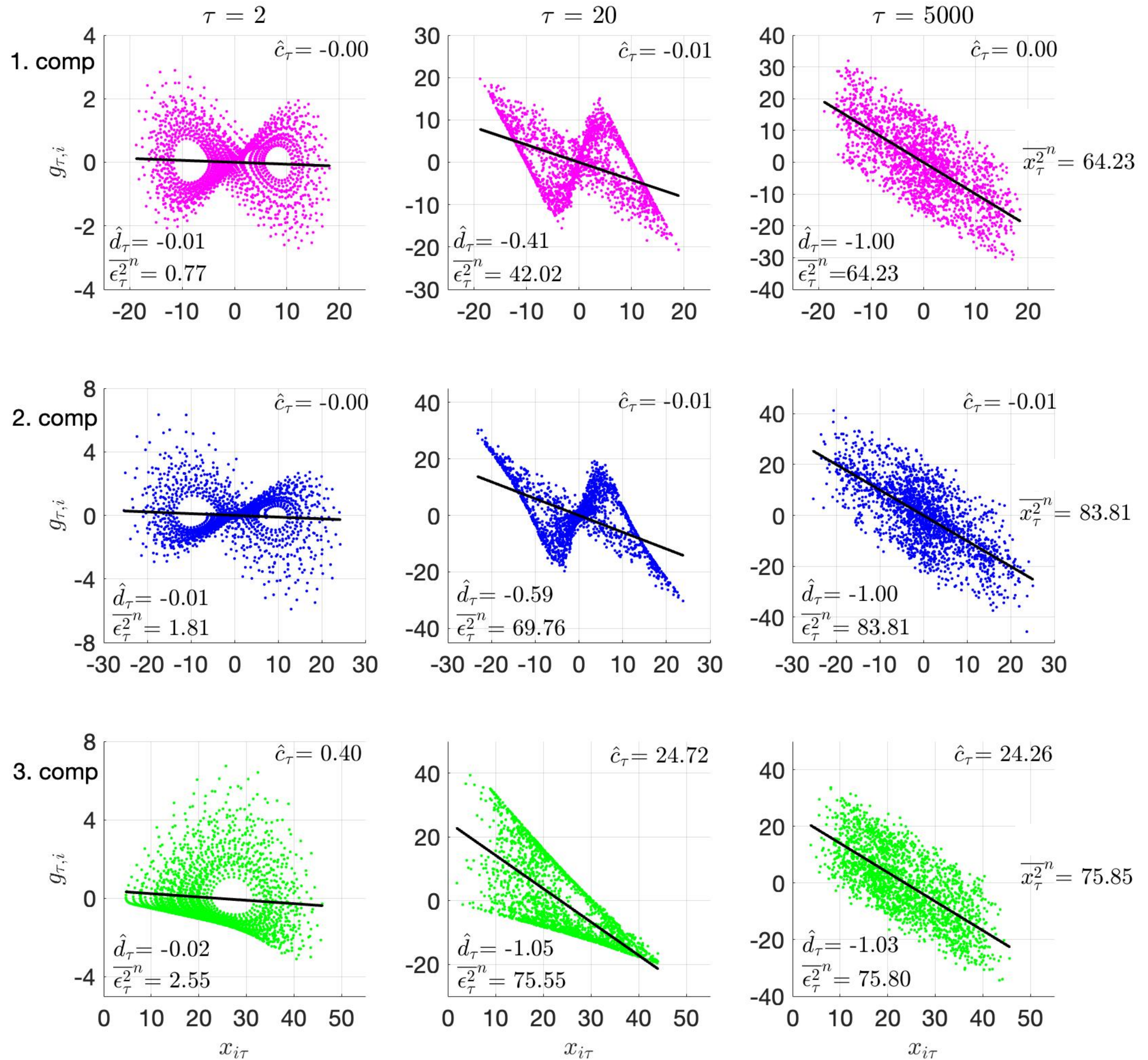
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$$g_{\tau,i} = \sum_{s=i\tau}^{(i+1)\tau-1} f_s = \hat{c}_\tau + \hat{d}_\tau x_{i\tau} + \hat{\epsilon}_{\tau,i}$$

dissipating component of  $g_{\tau,i}$  with strength  $|d_\tau|$

fluctuating component of  $g_{\tau,i}$  with strength  $\overline{\epsilon_\tau^2}^n$



- ◆  $g_{\tau,i}$  becomes increasingly linear in  $x_{i\tau}$  with increasing  $\tau$
- ◆ Once  $g_{\tau,i}$  is linear in  $x_{i\tau}$ ,  $\hat{\epsilon}_{\tau,i}$  behaves like a white noise
- ◆  $\hat{d}_\tau$  is always negative
- ◆  $\overline{\epsilon_\tau^2}^n$  increases with  $|\hat{d}_\tau|$ , and stops to increase and becomes equals to  $\overline{x_\tau^2}^n$  when  $|\hat{d}_\tau|=1$ , which happens when  $\tau > \tau_0$



# Properties of the integral forcing $g_\tau$

- The dissipating and fluctuating component of  $g_\tau$  are related to each other following the *FD-curve*:

$$\text{Var}(\epsilon_\tau) = \text{Var}(x) \left( 1 - (1 + d_\tau)^2 \right)$$

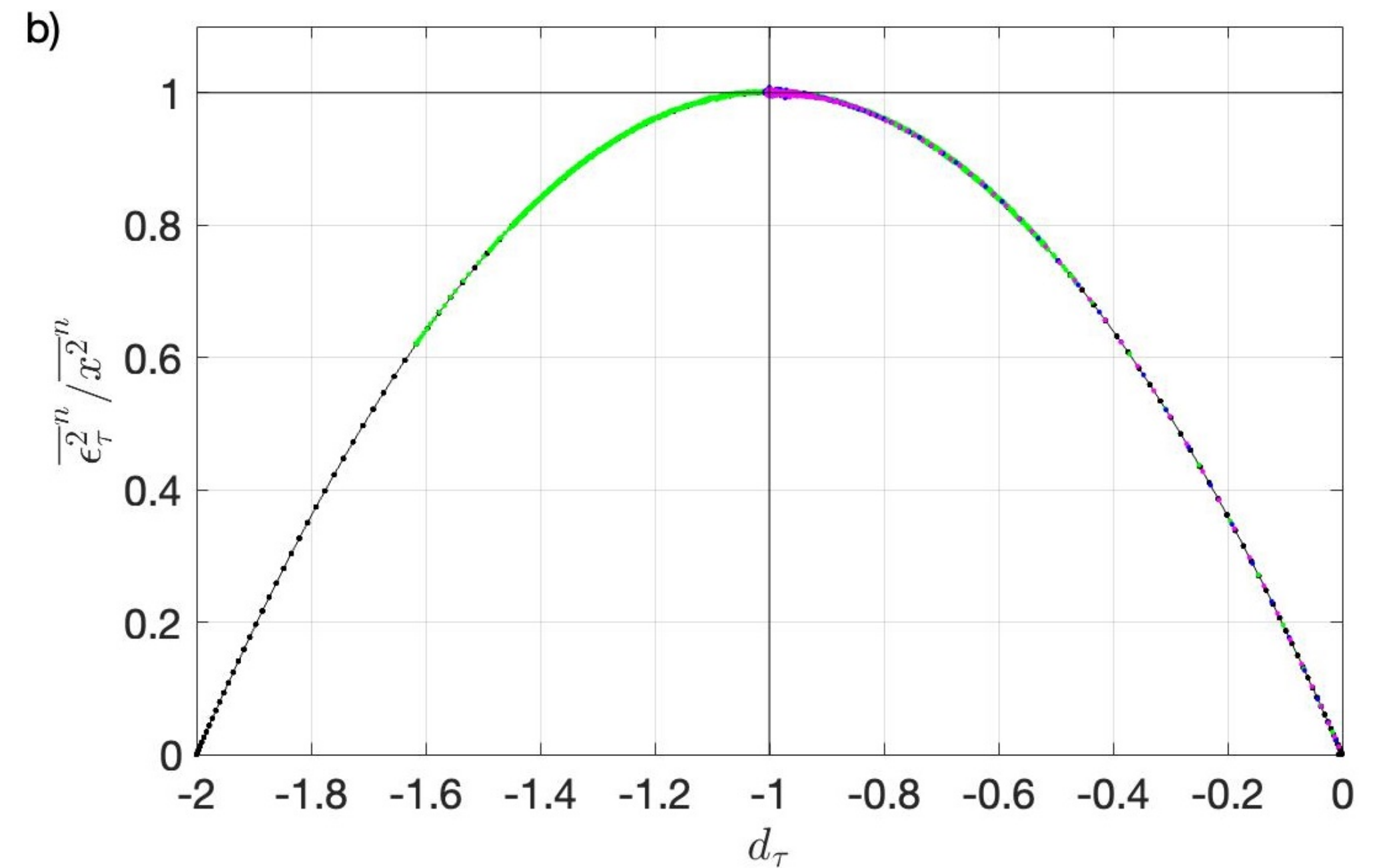
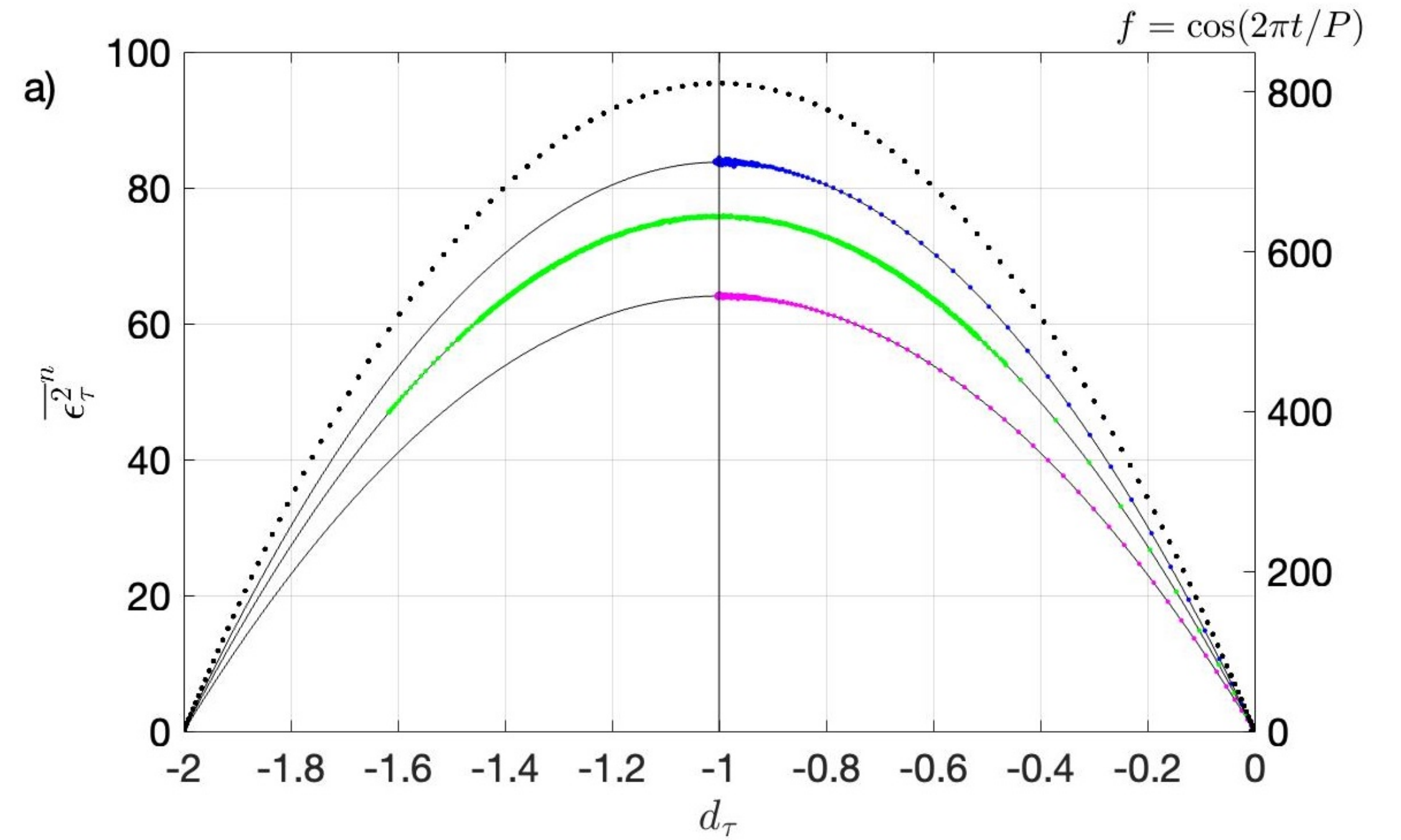
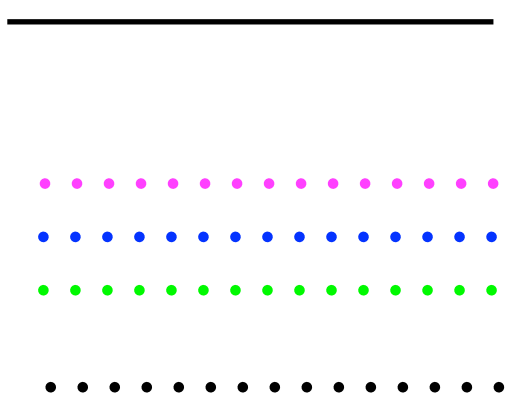
with  $\text{Var}(\epsilon_\tau) = \lim_{n \rightarrow \infty} \overline{\epsilon_\tau^{2n}}$ ,  $\text{Var}(x) = \lim_{n \rightarrow \infty} \overline{x^{2n}}$

- $(\hat{d}_\tau, \overline{\epsilon_\tau^{2n}})$ -points lie on the *DF-curve*
- $1 + d_\tau = \rho_\tau$  so that  $d_\tau \in [-2, 0]$
- $\overline{\epsilon_\tau^{2n}}$  reaches its maximum at  $d_\tau = -1$ , which equals  $\overline{x^{2n}}$
- The *FD-curve* is independent of the functional form of  $f$
- Different  $f$  make  $(d_\tau, \overline{\epsilon_\tau^{2n}})$ -points to populate different parts of the DF-curve

$y = a(1 - (1 + z)^2)$ :

$(\hat{d}_\tau, \overline{\epsilon_\tau^{2n}})$ ,  $\tau = 1, \dots, 10^3$  from Lorenz model:

$(\hat{d}_\tau, \overline{\epsilon_\tau^{2n}})$ ,  $\tau = 1, \dots, 10^3$  from  $dx/dt = \cos(2\pi t/P)$ :



# The integral effect = The ability of $g_\tau$ in producing white-noise like $\tau$ -stepping solution

- Even though equivalent in determining  $\{x_s\}$  at a time, the summation (needed for obtaining  $g_\tau$ ) makes  $f$  and  $g_\tau$  to contain different amounts of information about time sequence  
 →  $f$  and  $g_\tau$  generate variations in  $x$  in different ways

- $f_s$  generates a change  $x_{s+1} - x_s$   
 Variations in  $f$  at a frequency generate variations in  $x$  at the SAME frequency

following the classical forcing-response relation

- For  $g_\tau$  with  $\tau > \tau_0$ , we have  $x_{(i+1)\tau} = c_\tau + \epsilon_{\tau,i}$ , despite  $x_{(i+1)\tau} - x_{i\tau} = g_{\tau,i}$   
 $\{x_{i\tau}\}$  varies at all frequencies smaller than  $1/\tau_0$

Classical forcing-response relation does not exist:

- It is not meaningful to talk about a change  $x_{(i+1)\tau} - x_{i\tau}$ , since  $x_{(i+1)\tau}$  is independent of  $x_{i\tau}$
- It is also not meaningful to talk about  $g_{\tau,i}$  as a forcing, since  $x_{(i+1)\tau}$  equals only part of  $g_{\tau,i}$

Since  $x_{(i+1)\tau} = x_{i\tau} + g_{\tau,i} = x_{i\tau} + (c_\tau + d_\tau x_{i\tau} + \epsilon_{\tau,i})$ ,

$$x_{(i+1)\tau} = (1 + d_\tau)^{i+1} x_0 + \sum_{k=0}^i (1 + d_\tau)^k (c_\tau + \epsilon_{\tau,i-k})$$

~~$x_{(i+1)\tau} = x_0 + \sum_{k=0}^i (c_\tau + \epsilon_{\tau,i-k})$~~

$x_{(i+1)\tau} = c_\tau + \epsilon_{\tau,i}$

Like random walk:  
No equilibrium solution possible!

Like white noise!

→ Randomness emerges



■ **Integral effect does not exist for all integral forcing  $g_\tau$**

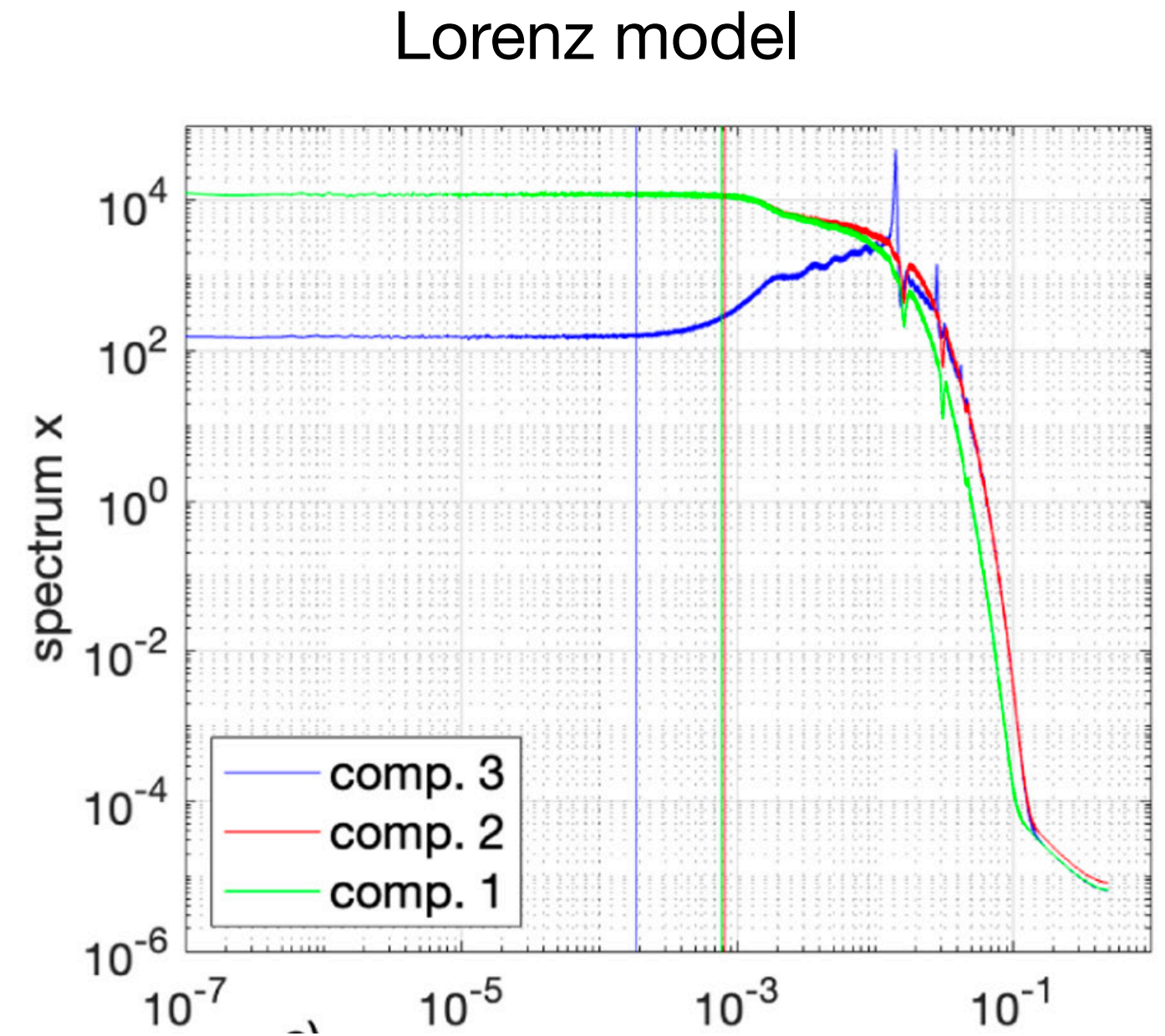
Since

- $f$  is unable to generate variations at the lowest frequencies
- $g_\tau$  with  $\tau > \tau_0$  exists only at frequencies  $\omega < 1/\tau_0$

the integral effect can be quantified in terms of  $Var(0) = 2\omega_0\Gamma(0)$  by

$$\frac{Var(0)}{Var(x)} = 2\omega_0 \left( 2 \sum_{\tau=1}^{\infty} (1 + d_\tau) + 1 \right)$$

where  $\Gamma(0) = \sum_{\tau=-\infty}^{\infty} \gamma_\tau$ ,  $[-\omega_0, \omega_0]$  with  $\omega_0 < 1/\tau_0$  is the frequency range over which the spectrum of  $x$  has a white extension



	$f^1$	$f^2$	$f^3$	$\cos(2\pi t/P)$
$2 \sum_{\tau=1}^n (1 + d_\tau) + 1$	86.45	68.13	1.86	-0.00
$\omega_0$	$7.03 \times 10^{-04}$	$7.05 \times 10^{-04}$	$7.20 \times 10^{-05}$	0
$Var(0)/Var(x) \times 100\%$	12.16%	9.61%	0.03%	0.00%

The integral effect is strongest for the first two Lorenz components, but does not exist for  $dx/dy = \cos(2\pi t/P)$

## CONCLUSIONS:

1.  $g_\tau$  with  $\tau \in \mathbb{Z}_+$  obeys a law-like regularity (FD-curve) that relates its dissipating component characterized by  $d_\tau$  to its fluctuating component characterized by  $Var(\epsilon_\tau)$ , independent of the functional form of  $f$

2. Randomness is an intrinsic feature of  $dx/dt = f$  that

- results from the joint working of the dissipating and the fluctuating component of  $g_\tau$
- only „visible“ by integrating  $dx/dt = f$  forward in time

3. With respect to the equilibrium variance of  $x$ , time is irreversible and has an arrow !

Thanks!