

1. Solve the partial differential equation

$$u_t + x u_x = 0 \quad \text{in } \mathbb{R} \times (0, \infty),$$

$$u = g \quad \text{on } \mathbb{R} \times \{t=0\},$$

where $u = u(x, t)$ and $g \in C^1(\mathbb{R})$.

Hint: Show that $z(s) = u(x(s), t+s)$ is constant if $x'(s) = x(s)$.

(10)

$$z'(s) = x'(s) u_x(x(s), t+s) + u_t(x(s), t+s)$$

Therefore, if $x'(s) = x(s)$,

$$z'(s) = x(s) u_x(x(s), t+s) + u_t(x(s), t+s)$$

$$= (x u_x + u_t)(x(s), t+s) = 0$$

$$\text{But } x'(s) = x(s) \Rightarrow x(s) = x_0 e^s$$

$$\text{Further: } z(0) = u(x_0, t)$$

$$z(-t) = u(x_0 e^{-t}, 0) = g(x_0 e^{-t})$$

Since $z(0) = z(-t)$, we obtain (writing x in place of x_0),

$$u(x, t) = g(x e^{-t}).$$

2. For $U \subset \mathbb{R}^n$ open and (path-)connected, let $u \in C^2(U)$ be harmonic with $u \geq 0$. Show that if $u(x) > 0$ for some $x \in U$, then $u > 0$ everywhere in U . (10)

Assume the contrary, i.e. $\exists y \in U$ s.t. $u(y) = 0$.

Let $\gamma: [0, 1] \rightarrow U$ denote a path connecting x and y , i.e.

$$\gamma(0) = x, \quad \gamma(1) = y$$

Then $\exists t^* \in (0, 1]$ s.t. $u(\gamma(t)) = 0$ for $t \in [t^*, 1]$ (*)

$u(\gamma(t)) > 0$ for $t < t^*$ near t^*

Let $x^* = \gamma(t^*)$. By the mean value formula,

$$u(x^*) = \int_{B(x^*, r)} u(x) dx$$

where r is small enough s.t. $B(x^*, r) \subset U$. Since u is continuous and $u > 0$ somewhere (along the path γ) in $B(x^*, r)$, the integral is strictly positive and $u(x^*) > 0$.

This contradicts (*).

3. Recall that the solution to the heat equation

$$\begin{aligned} u_t - \Delta u &= 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ u &= g & \text{on } \mathbb{R}^n \times \{t=0\} \end{aligned}$$

is given by

$$u(x, t) = \int_{\mathbb{R}^n} \Phi(x-y, t) g(y) dy,$$

where, for $t > 0$,

$$\Phi(z, t) = \frac{1}{(4\pi t)^{n/2}} e^{-\frac{|z|^2}{4t}}.$$

Assume that g is continuous and compactly supported. Show that there exists a $C > 0$, depending only on the support of g , such that

$$\|Du(x, t)\| \leq \frac{C}{\sqrt{t}} \|g\|_{L^\infty}. \quad (10)$$

$$|Du(x, t)| = \left| \int_{\mathbb{R}^n} \underbrace{D\Phi(x-y, t)}_{= \frac{1}{(4\pi t)^{n/2}} \left(-\frac{x-y}{2t}\right)} g(y) dy \right|$$

$$\leq \frac{1}{2^{n+1} \pi^{n/2} t^{n/2+1}} \int_{\mathbb{R}^n} |x-y| e^{-\frac{|x-y|^2}{4t}} dy \|g\|_{L^\infty}$$

$$= \int_{\mathbb{R}^n} \underbrace{|z| e^{-\frac{|z|^2}{4t}} dz}_{= \int_0^{\frac{r^2}{4t}} \int_{\partial B(0, \sqrt{s})} e^{-s} ds} = \int_0^{\frac{r^2}{4t}} n \omega(n) r^{n-1} dr$$

$$= n \omega(n) \int_0^{\frac{r^2}{4t}} s^{\frac{n}{2}-1} e^{-s} ds$$

$$= \frac{n \omega(n)}{\pi^{n/2}} \int_0^{\frac{r^2}{4t}} s^{\frac{n}{2}-1} ds \frac{1}{\sqrt{t}} \|g\|_{L^\infty}$$

$$= C$$

$$= \frac{n \omega(n)}{\pi^{n/2}} \int_0^{\frac{r^2}{4t}} s^{\frac{n}{2}-1} ds \frac{1}{\sqrt{t}} \|g\|_{L^\infty}$$

$$= C$$

4. Let $U \subset \mathbb{R}^n$ be open and bounded with C^1 boundary. Assume that $u \in C^2(\bar{U} \times [0, \infty))$ solves the heat equation

$$\begin{aligned} u_t - \Delta u &= 0 & \text{in } U \times (0, \infty), \\ u &= g & \text{on } U \times \{t=0\}, \\ u &= 0 & \text{on } \partial U \times (0, \infty), \end{aligned}$$

and define the "energy"

$$E(t) = \int_U |u(x, t)|^2 dx.$$

Prove that $E(t) \leq E(0)$ for every $t \geq 0$.

(10)

Multiply the heat equation with u and integrate in space:

$$\int_U u_t u dx = \int_U \Delta u u dx$$

$$\Rightarrow \frac{1}{2} \frac{d}{dt} \int_U |u|^2 dx = \underbrace{\int_U u \nabla \cdot Du ds}_{=0} - \underbrace{\int_U |Du|^2 dx}_{\leq 0}$$

$\Rightarrow E(t)$ is non-increasing.

$$y = x + tz$$

$$\Rightarrow ds(y) = t^2 ds(z)$$

5. Let $h \in C^2(\mathbb{R}^3)$ and set

$$u(x, t) = \int_{\partial B(x, t)} t h(y) ds(y)$$

for $t > 0$.

Show that

$$= t \frac{1}{3\alpha(3)} t^2 \int_{\partial B(x, t)} h(y) ds(y)$$

(a) $\lim_{t \rightarrow 0} u(x, t) = 0$.

(b) $\lim_{t \rightarrow 0} u_t(x, t) = h(x)$.

(c) (Extra credit.) u solves the wave equation $\Delta u = 0$

$$u_{tt} - \Delta u = 0$$

on $\mathbb{R}^3 \times (0, \infty)$.

$$(5+5+5)$$

(a) $u(x, t) = t \int_{\partial B(x, t)} h(y) ds(y) \rightarrow 0$ as $t \rightarrow 0$

$\rightarrow h(x)$ as $t \rightarrow 0$

(b) $u_t = \int_{\partial B(x, t)} h(y) ds(y) + \frac{t}{3\alpha(3)} \int_{\partial B(0, 1)} z \cdot \nabla h(x+tz) ds(z)$

$$= \frac{1}{t^2} \int_{\partial B(x, t)} z \cdot \nabla h(y) ds(y)$$

$$= \frac{1}{t^2} \int_{\partial B(x, t)} \Delta h dy$$

$$= \int_{\partial B(x, t)} h(y) ds(y) + \frac{t^2}{3} \int_{\partial B(x, t)} \Delta h dy$$

$\rightarrow h(x)$ $\rightarrow 0$

(c) $u_t = \frac{1}{3\alpha(3)} t^2 \int_{\partial B(x, t)} h(y) ds(y) + \frac{1}{3\alpha(3)} t \int_0^t \Delta h ds(y) ds$

$$= \frac{1}{3\alpha(3)} \int_{\partial B(0, 1)} h(x+tz) ds(z)$$

$$\Rightarrow u_{tt} = \frac{1}{3\alpha(3)} \int_{\partial B(0, 1)} z \cdot \nabla h(x+tz) ds(z) - \frac{1}{3\alpha(3)} t^2 \int_{\partial B(x, t)} \Delta h dy$$

$$+ \frac{1}{3\alpha(3)} t \int_{\partial B(x, t)} \Delta h ds(y)$$

Changing back from z to y in the first integral, and applying the Gauss integral formula as in (b), we see that the first two terms cancel.

$$\Rightarrow u_{tt} = \frac{1}{3\alpha(3)} t \int_{\partial B(x, t)} \Delta ds(y)$$

On the other hand, taking Δ of (*),

$$\Delta u = \frac{t}{3\alpha(3)} \int_{\partial B(0, 1)} \Delta h(x+tz) ds(z)$$

$$= \frac{t}{3\alpha(3)} t^2 \int_{\partial B(x, t)} \Delta h ds(y)$$

$$\Rightarrow u_{tt} = \Delta u$$

□